

Dynamics and homogenised properties of disordered cellular structures

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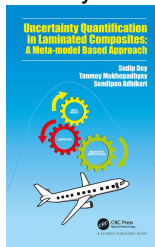
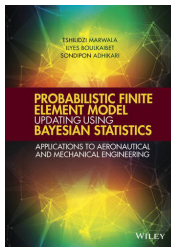
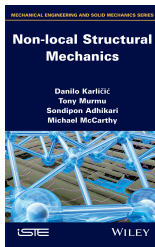
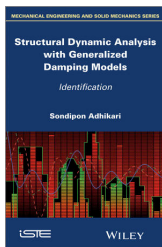
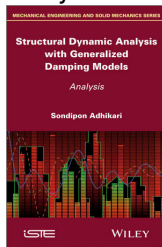
National University of Defense Technology, Changsha, 18 December 2019







- *Development* of fundamental computational methods for structural dynamics and uncertainty quantification
 - A. Dynamics of complex systems
 - B. Inverse problems for linear and non-linear dynamics
 - C. Uncertainty quantification in computational mechanics
- *Applications* of computational mechanics to emerging multidisciplinary research areas
 - D. Vibration energy harvesting / dynamics of wind turbines
 - E. Dynamics and mechanics of metamaterials and multi-scale systems



- 1 Introduction**
 - A brief history of metamaterials
 - Regular lattices
 - Irregular lattices

- 2 Formulation for the viscoelastic analysis**

- 3 Equivalent elastic properties of randomly irregular lattices**

- 4 Effective properties of irregular lattices: uncorrelated uncertainty**

- General results - closed-form expressions
- Special case 1: Only spatial variation of the material properties
- Special case 2: Only geometric irregularities
- Special case 3: Regular hexagonal lattices

- 5 Effective properties of irregular lattices: correlated uncertainty**

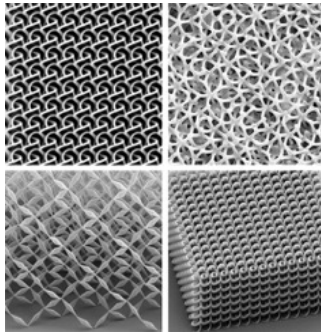
- 6 Results and discussions**

- Spatially correlated irregular elastic lattices
- Viscoelastic properties of regular lattices
- Spatially correlated irregular viscoelastic lattices

- 7 Frequency-dependent effective elastic moduli**

- 8 Conclusions**

- Metamaterials are **designer media** with periodic units comprised of unique tailor-made geometry and pattern aimed at accomplishing exceptional and unusual bulk properties which are unprecedented in conventional materials.
- Mechanical metamaterials achieve their **unusual effective properties** from the geometry and structure, and not from the intrinsic property of the constitutive material.
- An essential feature of metamaterials, operating in any frequency ranges or of any length-scales, is the **periodicity of a unit cell**:



- The origin of metamaterials was in field of *electromagnetism* and early ideas can be traced back to 1967 [1–3].
- However, it was the seminal paper by Smith et al. [4] in 2000 demonstrating **negative** permeability and permittivity by a periodic array of split-ring resonators, that started the current interest in metamaterials.
- Since then several concepts and devices have been conceived which **challenge conventional physical laws**, such as negative refraction [5], the perfect lens [6, 7], and invisibility cloaking in electromagnetism and optics [8, 9].
- These **extraordinary developments** not only attracted researchers but also captured the imagination of the public and in some cases, science fiction [10].

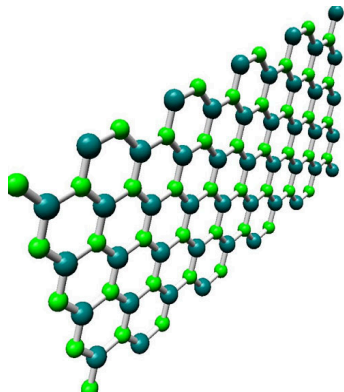
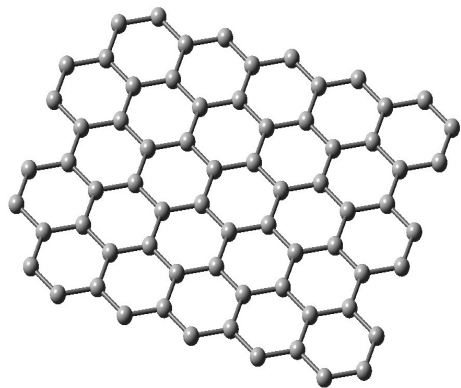
- The next round of development was in *acoustic* metamaterials [11, 12] exploiting the idea of locally resonant behaviour [13].
- The consideration of sub-wavelength structures [14] opened up immense possibilities, including *negative* effective elastic modulus [15], *negative* density (or mass) [16], or both [17], anisotropy in the effective mass or density [18, 19], and *non-reciprocal* response [20, 21].
- The *mechanical* metamaterials emerged following in the footsteps of electromagnetic [22] and acoustic metamaterials [23], primarily within the past five years [24–28].
- *Intense research* in recent years shows ultralight metamaterials [29] approaching theoretical strength limit [30], pentamode materials [31] with cloaking mode [32], negative refraction elastic waves [33], elastic cloaking [34, 35] and hyperbolic elastic metamaterials [36].

- The **rise of mechanical metamaterials** [26, 28] coincides with remarkable recent advances in manufacturing technology [37].
- From the point of view of analytical techniques, **two distinct type** of metamaterials are considered, namely, (1) static mechanical metamaterials and (2) dynamic mechanical metamaterials.
- The roots of static mechanical metamaterials can be traced back to late 80's with the discovery of **negative Poisson's ratio** cellular structures [38].
- This talk is related to pseudo-dynamic behaviour of **viscoelastic** cellular metamaterials with disordered geometry and material properties.
- The viscoelastic behaviour makes the equivalent in-plane properties **frequency dependent**.

- **Lattice** based metamaterials are abundant in man-made and natural systems at various length scales.
- They are made of **periodic** identical/near-identical geometric units.
- Among various lattice geometries (triangle, square, rectangle, pentagon, hexagon), **hexagonal lattice** is most common (note that hexagon is the highest “space filling” pattern in 2D).
- This talk is about **in-plane** viscoelastic properties of 2D hexagonal lattice structures - commonly known as “honeycombs”

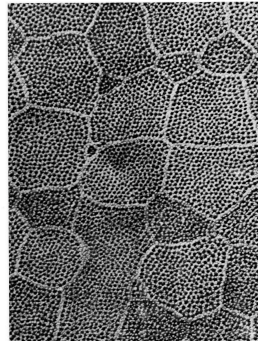
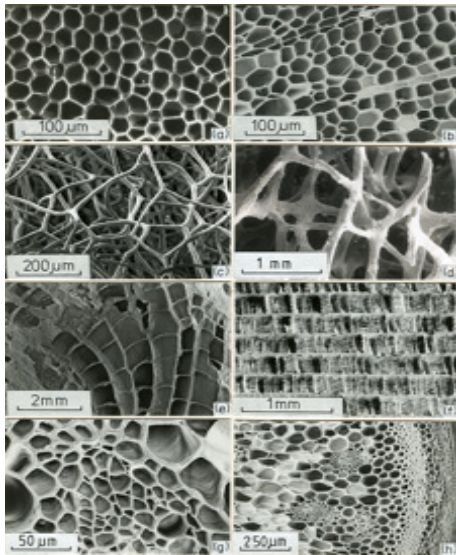


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Illustrations of a single layer graphene sheet and a boron nitride nano sheet

Lattice structures - nature

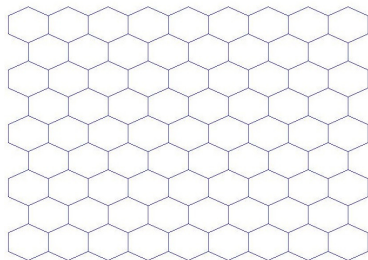


Top left: cork, top right: balsa, next down left: sponge, next down right: trabecular bone, next down left: coral, next down right: cuttlefish bone, bottom left: leaf tissue, bottom right: plant stem, third column - epidermal cells (from web.mit.edu)

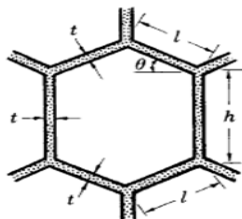
- Shall we consider lattices as “structures” or “materials” from a mechanics point of view?
- At what relative length-scale a lattice *structure* can be considered as a *material* with equivalent elastic properties?
- In what ways structural irregularities *influence* equivalent elastic / viscoelastic properties? Can we evaluate it in a quantitative as well as in a qualitative manner?
- What is the consequence of *random* structural irregularities on the homogenisation approach in general? Can we obtain statistical measures?
- How can we efficiently *compute* equivalent elastic / viscoelastic properties of random lattice structures?

- Hexagonal lattice structures have been modelled as a **continuous solid** with an equivalent elastic moduli throughout its domain.
- This approach **eliminates** the need of detail finite element modelling of lattices in complex structural systems like sandwich structures.
- Extensive amount of research has been carried out to predict the **equivalent elastic / viscoelastic properties** of regular lattices consisting of perfectly periodic hexagonal cells [39].
- Analysis of two dimensional hexagonal lattices dealing with **in-plane elastic properties** are commonly based on an unit cell approach, which is applicable only for perfectly periodic cellular structures.
- For the dynamic analysis of perfectly periodic structures, Floquet-Bloch theorem is normally employed to characterise wave propagation.

- Unit cell approach - Gibson and Ashby (1999)



(a) Regular hexagon ($\theta = 30^\circ$)



(b) Unit cell

- We are interested in homogenised equivalent in-plane elastic properties
- This way, we can avoid a detailed structural analysis considering all the beams and treat it as a material

- The cell walls are treated as beams of thickness t , depth b and Young's modulus E_s . l and h are the lengths of inclined cell walls having inclination angle θ and the vertical cell walls respectively.
- The equivalent elastic properties are:

$$E_1 = E_s \left(\frac{t}{l}\right)^3 \frac{\cos \theta}{\left(\frac{h}{l} + \sin \theta\right) \sin^2 \theta} \quad (1)$$

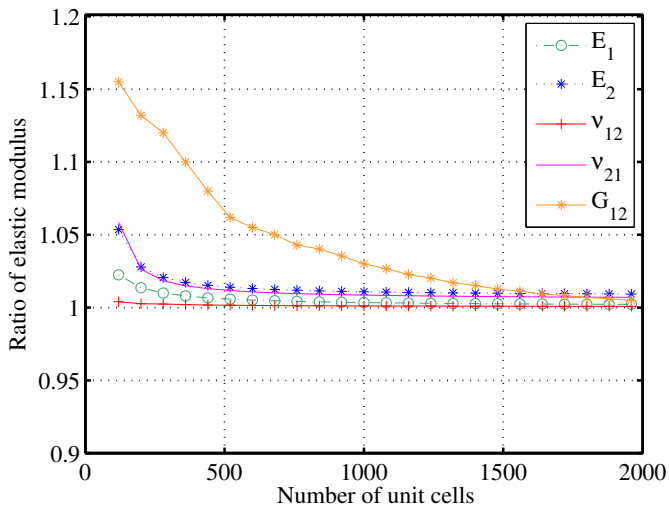
$$E_2 = E_s \left(\frac{t}{l}\right)^3 \frac{\left(\frac{h}{l} + \sin \theta\right)}{\cos^3 \theta} \quad (2)$$

$$\nu_{12} = \frac{\cos^2 \theta}{\left(\frac{h}{l} + \sin \theta\right) \sin \theta} \quad (3)$$

$$\nu_{21} = \frac{\left(\frac{h}{l} + \sin \theta\right) \sin \theta}{\cos^2 \theta} \quad (4)$$

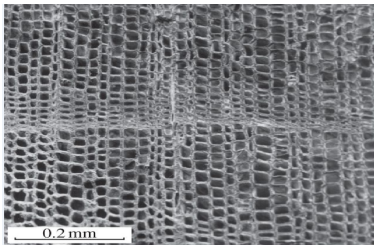
$$G_{12} = E_s \left(\frac{t}{l}\right)^3 \frac{\left(\frac{h}{l} + \sin \theta\right)}{\left(\frac{h}{l}\right)^2 \left(1 + 2\frac{h}{l}\right) \cos \theta} \quad (5)$$

- A finite element code has been developed to obtain the in-plane elastic moduli numerically for hexagonal lattices.
- Each cell wall has been modelled as an Euler-Bernoulli beam element having three degrees of freedom at each node.
- For E_1 and ν_{12} : two opposite edges parallel to direction-2 of the entire hexagonal lattice structure are considered. Along one of these two edges, uniform stress parallel to direction-1 is applied while the opposite edge is restrained against translation in direction-1. Remaining two edges (parallel to direction-1) are kept free.
- For E_2 and ν_{21} : two opposite edges parallel to direction-1 of the entire hexagonal lattice structure are considered. Along one of these two edges, uniform stress parallel to direction-2 is applied while the opposite edge is restrained against translation in direction-2. Remaining two edges (parallel to direction-2) are kept free.
- For G_{12} : uniform shear stress is applied along one edge keeping the opposite edge restrained against translation in direction-1 and 2, while the remaining two edges are kept free.



$\theta = 30^\circ$, $h/l = 1.5$. FE results converge to analytical predictions after 1681 cells.

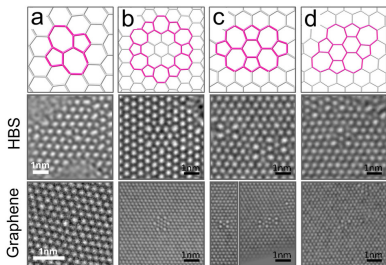
Irregular lattice structures



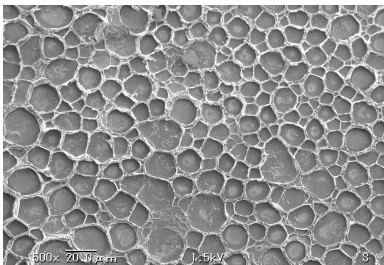
(c) Cedar wood



(d) Manufactured honeycomb core



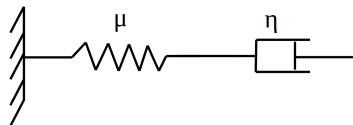
(e) Graphene image



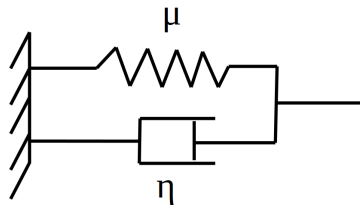
(f) Fabricated CNT surface

- A **significant limitation** of the aforementioned unit cell approach is that it cannot account for the spatial irregularity, which is practically inevitable.
- **Spatial irregularity** may occur due to manufacturing uncertainty, structural defects, variation in temperature, pre-stressing and micro-structural variabilities.
- To include the effect of irregularity, **voronoi honeycombs** have been considered in several studies.
- The effect of different forms of irregularity on elastic properties and structural responses of hexagonal lattices are generally based on **direct finite element (FE) simulation**.
- In the FE approach, a small change in geometry of a single cell may require completely new geometry and meshing of the entire structure. In general this makes the entire process **time consuming and tedious**.
- The problem becomes worse for **uncertainty quantification** of the responses, where the expensive finite element model is needed to be simulated for a large number of samples while using a Monte Carlo based approach.

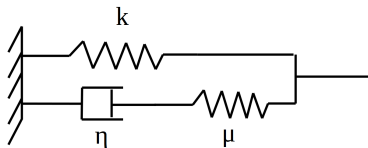
Viscoelastic models



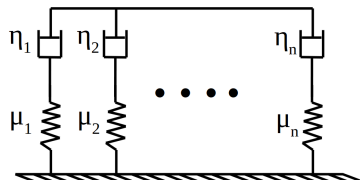
(g) Maxwell model



(h) Voigt model



(i) Standard linear model



(j) Generalised Maxwell model

Figure: Springs and dashpots based models viscoelastic materials.

The viscoelastic kernel function can be expressed for the four models as

- *Maxwell model:*

$$g(t) = \mu e^{-(\mu/\eta)t} \mathcal{U}(t) \quad (6)$$

- *Voigt model:*

$$g(t) = \eta \delta(t) + \mu \mathcal{U}(t) \quad (7)$$

- *Standard linear model:*

$$g(t) = E_R \left[1 - \left(1 - \frac{\tau_\sigma}{\tau_\epsilon} \right) e^{-t/\tau_\epsilon} \right] \mathcal{U}(t) \quad (8)$$

- *Generalised Maxwell model:*

$$g(t) = \left[\sum_{j=1}^n \mu_j e^{-(\mu_j/\eta_j)t} \right] \mathcal{U}(t) \quad (9)$$

Models similar to this is also known as the Pony series model.

Viscoelastic model	Complex modulus
Biot model	$G(\omega) = G_0 + \sum_{k=1}^n \frac{a_k i \omega}{i \omega + b_k}$
Fractional derivative	$G(\omega) = \frac{G_0 + G_\infty (i \omega \tau)^\beta}{1 + (i \omega \tau)^\beta}$
GHM	$G(\omega) = G_0 \left[1 + \sum_k \alpha_k \frac{-\omega^2 + 2i \xi_k \omega_k \omega}{-\omega^2 + 2i \xi_k \omega_k \omega + \omega_k^2} \right]$
ADF	$G(\omega) = G_0 \left[1 + \sum_{k=1}^n \Delta_k \frac{\omega^2 + i \omega \Omega_k}{\omega^2 + \Omega_k^2} \right]$
Step-function	$G(\omega) = G_0 \left[1 + \eta \frac{1 - e^{-st_0}}{st_0} \right]$
Half cosine model	$G(\omega) = G_0 \left[1 + \eta \frac{1 + 2(st_0/\pi)^2 - e^{-st_0}}{1 + 2(st_0/\pi)^2} \right]$
Gaussian model	$G(\omega) = G_0 \left[1 + \eta e^{\omega^2/4\mu} \left\{ 1 - \operatorname{erf} \left(\frac{i\omega}{2\sqrt{\mu}} \right) \right\} \right]$

Complex modulus for some viscoelastic models in the frequency domain

- We consider that each constitutive element of a hexagonal unit within the lattice structure is modelled using viscoelastic properties. For simplicity, we use Biot model with only one term. Frequency dependent complex elastic modulus for an element is expressed as

$$E(\omega) = E_S \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \quad (10)$$

where μ and ϵ are the relaxation parameter and a constant defining the 'strength' of viscosity, respectively. E_S is the intrinsic Young's modulus.

- The amplitude of this complex elastic modulus is given by

$$|E(\omega)| = E_S \sqrt{\frac{\mu^2 + \omega^2 (1 + \epsilon)^2}{\mu^2 + \omega^2}} \quad (11)$$

- The phase (ϕ) of this complex elastic modulus is given by

$$\phi(E(\omega)) = \tan^{-1} \left(\frac{\epsilon\mu\omega}{\mu^2 + \omega^2(1 + \epsilon)} \right) \quad (12)$$

Mathematical idealisation of irregularity in lattice structures

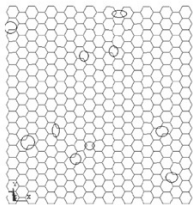


Fig. Randomly missing cell wall

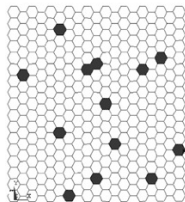


Fig. Random filled cell

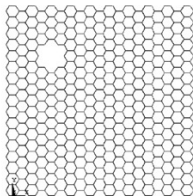
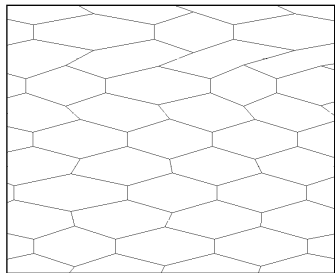
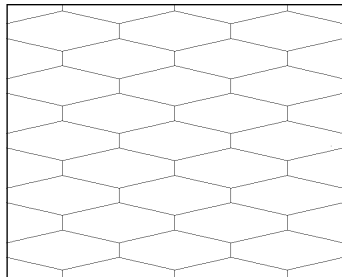


Fig. Missing cell cluster



- Random spatial irregularity in cell angle is considered in this study.

- The equivalent elastic properties for a regular lattice:

$$E_1 = E_s \left(\frac{t}{l} \right)^3 \frac{\cos \theta}{\left(\frac{h}{l} + \sin \theta \right) \sin^2 \theta} \quad (13)$$

$$E_2 = E_s \left(\frac{t}{l} \right)^3 \frac{\left(\frac{h}{l} + \sin \theta \right)}{\cos^3 \theta} \quad (14)$$

$$\nu_{12} = \frac{\cos^2 \theta}{\left(\frac{h}{l} + \sin \theta \right) \sin \theta} \quad (15)$$

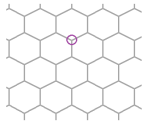
$$\nu_{21} = \frac{\left(\frac{h}{l} + \sin \theta \right) \sin \theta}{\cos^2 \theta} \quad (16)$$

$$G_{12} = E_s \left(\frac{t}{l} \right)^3 \frac{\left(\frac{h}{l} + \sin \theta \right)}{\left(\frac{h}{l} \right)^2 \left(1 + 2 \frac{h}{l} \right) \cos \theta} \quad (17)$$

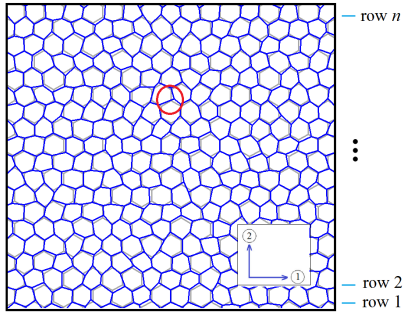
- Parameters of these expressions **CANNOT** be randomised to obtain equivalent properties for an irregular lattice.

- Direct numerical simulation to deal with irregularity in lattice structures may not necessarily provide proper understanding of the underlying physics of the system. An **analytical approach** could be a simple, insightful, yet an efficient way to obtain effective elastic properties of lattice structures.
- This work develops a structural mechanics based analytical framework for predicting equivalent in-plane elastic properties of irregular lattices having **spatially random** variations in cell angles.
- **Closed-form** analytical expressions will be derived for equivalent in-plane elastic properties.
- An approach based on the **frequency-domain** representation of the viscoelastic property of the constituent elements in the cells is used.

The philosophy of the analytical approach for irregular lattices



(a)

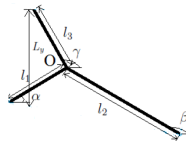
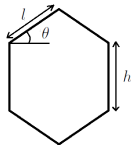


— regular honeycomb — irregular honeycomb

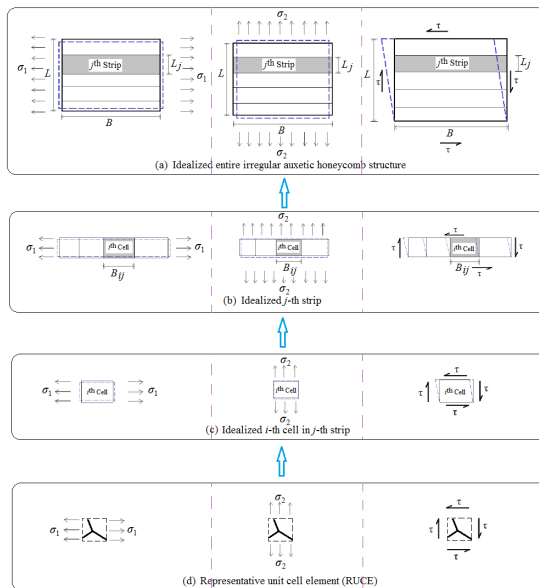
(b)

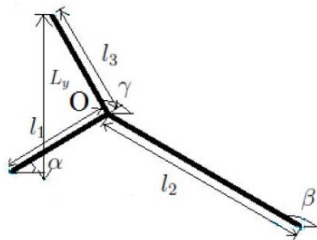
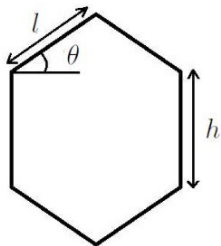


(c)



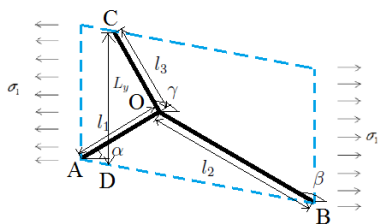
The idealisation of RUCE and the bottom-up homogenisation approach



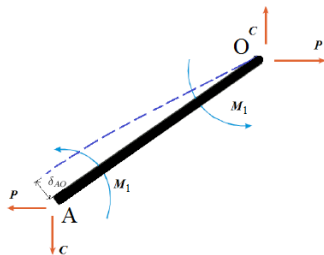


(a) Classical unit cell for regular lattices (b) Representative unit cell element (RUCE) geometry for irregular lattices

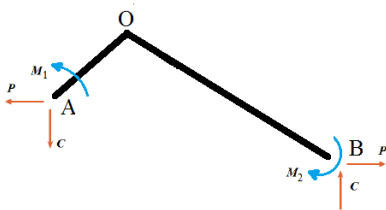
RUCE and free-body diagram for the derivation of E_1



(a)



(b)



(c)

Longitudinal Young's modulus for an idealized RUCE

- Stress σ_1 is applied in direction-1 for deriving the expression of longitudinal Young's modulus for a single RUCE (E_{1U}). From the condition of vertical equilibrium, it can be concluded that the vertical forces acting on points A and B should be of equal magnitude and opposite sign.
- The horizontal forces acting on points A and B can be expressed as $P = \sigma_1 L_y b$, where L_y represents the length CD and b is the height of honeycomb sheet (dimension perpendicular to the 1-2 plane).
- The moments M_1 and M_2 can be expressed as

$$M_1 = \frac{1}{2}(Pl_1 \sin \alpha - Cl_1 \cos \alpha) \quad (18)$$

$$M_2 = \frac{1}{2}(Pl_2 \sin \beta - Cl_2 \cos \beta) \quad (19)$$

- Considering the rotational equilibrium of the free-body diagram presented in, the expression for C can be obtained as

$$C = P \left(\frac{l_1 \sin \alpha - l_2 \sin \beta}{l_1 \cos \alpha - l_2 \cos \beta} \right) \quad (20)$$

- The horizontal deflection of point A with respect to point O (δ_{AO}^h) consists of the deflection due to force P and the force C

$$\delta_{AO}^h = \left(\frac{Pl_1^3 \sin \alpha}{12E_s I} - \frac{Cl_1^3 \cos \alpha}{12E_s I} \right) \sin \alpha \quad (21)$$

where the first and second terms in the bracket represents the deflection of point A with respect to point O in the direction perpendicular to AO due to forces P and C respectively.

- The superscript h is used to represent horizontal direction of the applied stress. Here, E_s represents the intrinsic material property of the material, by which the honeycomb cell walls (/connecting members) are made of.
- The notation I represents the second moment of area of the cell walls, i.e. $I = bt^3/12$, where t denotes the thickness of honeycomb cell wall.

- The horizontal deflection of point B with respect to point O can be expressed as

$$\delta_{BO}^h = \left(\frac{Pl_2^3 \sin \beta}{12E_s I} - \frac{Cl_1^3 \cos \beta}{12E_s I} \right) \sin \beta \quad (22)$$

- The distance of the point vertically below joint O and on the line AB is given by

$$\delta_O = \frac{l_2 \sin \beta l_1 \cos \alpha - l_1 \sin \alpha l_2 \cos \beta}{l_1 \cos \alpha - l_2 \cos \beta} \quad (23)$$

Considering a linear strain field along the line AB, the effective horizontal deformation of the RUCE is given by

$$\begin{aligned} \delta_1^h &= \delta_{AO}^h \frac{\delta_O}{l_1 \sin \alpha} + \delta_{BO}^h \frac{\delta_O}{l_2 \sin \beta} \\ &= \frac{\sigma_1 L_y l_1^2 l_2^2 (l_1 + l_2) (\cos \alpha \sin \beta - \sin \alpha \cos \beta)^2}{E_s t^3 (l_1 \cos \alpha - l_2 \cos \beta)^2} \end{aligned} \quad (24)$$

- The strain in direction-1 can be obtained from 24 as

$$\epsilon_1^h = \frac{\sigma_1 L_y l_1^2 l_2^2 (l_1 + l_2) (\cos \alpha \sin \beta - \sin \alpha \cos \beta)^2}{E_s t^3 (l_1 \cos \alpha - l_2 \cos \beta)^3} \quad (25)$$

From 25, elastic modulus of a single RUCE in direction-1 is expressed as

$$E_{1U} = \frac{E_s t^3 (l_1 \cos \alpha - l_2 \cos \beta)^3}{L_y l_1^2 l_2^2 (l_1 + l_2) (\cos \alpha \sin \beta - \sin \alpha \cos \beta)^2} \quad (26)$$

Longitudinal Young's modulus for a Non-idealized RUCE

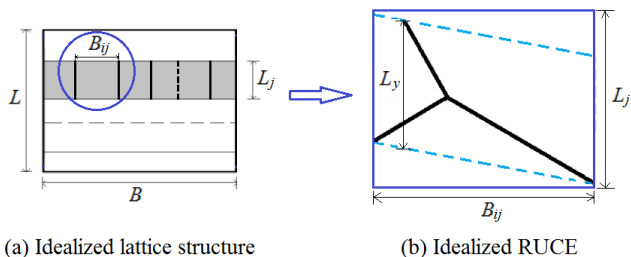


Figure: Idealization scheme of RUCE and the irregular lattice structure

- The expression of E_{1U} is for a non-idealized RUCE having a dimension of L_y in direction-2. However, for assembling the local properties of RUCes conveniently to the global level, it is essential to obtain the equivalent material property of an idealized RUCE (E'_{1U}) that has a virtual dimension of L_j (dimension of the j^{th} strip in direction-2).

- Considering a linear strain field, E'_{1U} can be obtained based on the deformation compatibility condition along direction-1, i. e. the deformation of the idealized RUCE and non-idealized RUCE in direction-1 should be equal

$$\frac{PB_{ij}}{A_{NI}E_{1U}} = \frac{PB_{ij}}{A_I E'_{1U}} \quad (27)$$

Here $A_{NI} = L_y b$ and $A_I = L_j b$. The above equation can be reduced to

$$E'_{1U} = E_{1U} \frac{L_y}{L_j} \quad (28)$$

- The deformation compatibility of j^{th} strip ensures that the total deformation of the strip in direction-1 due to stress σ_1 (Δ_{1j}) is the summation of individual deformations in direction-1 of each idealized RUC (Δ_{1ij}), while deformation of the idealized RUCs of that strip in direction-2 are same. Thus for the j^{th} strip

$$\Delta_{1j} = \sum_{i=1}^m \Delta_{1ij} \quad (29)$$

- The 29 can be rewritten as

$$\epsilon_{1j} B_j = \sum_{i=1}^m \epsilon_{1ij} B_{ij} \quad (30)$$

where ϵ_{1j} and B_j represent total strain and dimension in direction-1 for the j^{th} strip. Here $B_{ij} = (l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})$ and $B_j = \sum_{i=1}^m B_{ij}$.

- Equation (30) leads to

$$\frac{\sigma_1 B_j}{\hat{E}_{1j}} = \sum_{i=1}^m \frac{\sigma_1 B_{ij}}{E'_{1Uij}} \quad (31)$$

From 31, equivalent Young's modulus of j^{th} strip (\hat{E}_{1j}) can be expressed as

$$\hat{E}_{1j} = \frac{B_j}{\sum_{i=1}^m \frac{B_{ij}}{E'_{1Uij}}} \quad (32)$$

where E'_{1Uij} is the equivalent longitudinal elastic modulus in direction-1 of a single idealized RUCE positioned at (i,j) that can be obtained from equation (28).

- In the next step, closed-form expression for equivalent longitudinal Young's modulus of the entire irregular lattice (E_{1eq}) is obtained using the equivalent longitudinal Young's modulus for a single strip (\hat{E}_{1j}).

- Employing the force equilibrium conditions and deformation compatibility condition we have

$$\sigma_1 L b = \sum_{j=1}^n \sigma_{1j} L_j b \quad (33)$$

where L_j is the dimension of j^{th} strip in direction-2 and $L = \sum_{j=1}^n L_j$. The notation b represents the dimension of the lattice in the perpendicular direction to 1-2 plane.

- As strains in direction-1 for each of the n strips are the same to satisfy the deformation compatibility condition, equation (33) leads to

$$E_{1eq} L = \sum_{j=1}^n \hat{E}_{1j} L_j \quad (34)$$

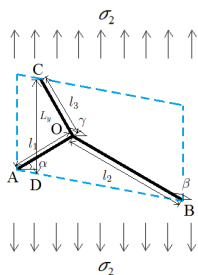
- Using 32 and 34, the equivalent Young's modulus in direction-1 of the entire irregular honeycomb structure (E_{1eq}) can be expressed as

$$E_{1eq} = \frac{1}{L} \sum_{j=1}^n \frac{B_j L_j}{\sum_{i=1}^m \frac{B_{ij}}{E_{1Uij}}} \quad (35)$$

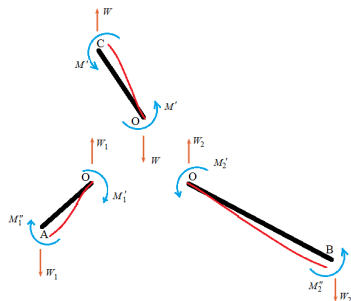
- From equations (26), (28) and (35), the expression for the longitudinal elastic modulus of the entire irregular lattice can be written as

$$E_{1eq} = \frac{E_s t^3}{L} \sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\sum_{i=1}^m \frac{l_{1ij}^2 l_{2ij}^2 (l_{1ij} + l_{2ij}) (\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})^2}{(l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})^2}} \quad (36)$$

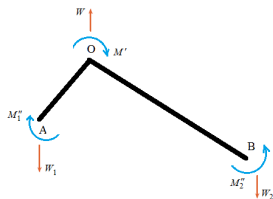
RUCE and free-body diagram for the derivation of E_2



(a)

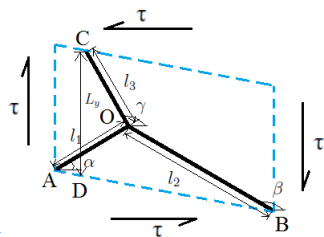


(b)

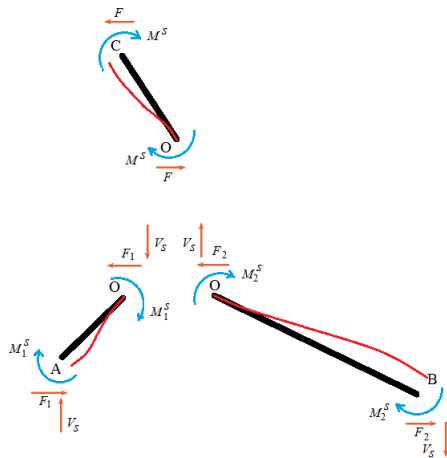


(c)

RUCE and free-body diagram for the derivation of G_{12}



(a)



(b)

Equivalent E_1

$$E_{1v}(\omega) = \frac{t^3}{L} \sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\sum_{i=1}^m \frac{l_{1ij}^2 l_{2ij}^2 (l_{1ij} + l_{2ij}) (\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})^2}{E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega}\right) ((l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})^2)}} \quad (37)$$

 Equivalent Young's moduli E_2

$$E_{2v}(\omega) = \frac{Lt^3}{\sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\sum_{i=1}^m E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega}\right) \left(l_{3ij}^2 \cos^2 \gamma_{ij} \left(l_{3ij} + \frac{l_{1ij} l_{2ij}}{l_{1ij} + l_{2ij}} \right) + \frac{l_{1ij}^2 l_{2ij}^2 (l_{1ij} + l_{2ij}) \cos^2 \alpha_{ij} \cos^2 \beta_{ij}}{(l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})^2} \right)^{-1}}} \quad (38)$$

Equivalent G_{12}

$$G_{12v}(\omega) = \frac{Lt^3}{\sum_{j=1}^n \frac{m}{\sum_{i=1}^m E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega} \right) \left(l_{3ij}^2 \sin^2 \gamma_{ij} \left(l_{3ij} + \frac{l_{1ij}l_{2ij}}{l_{1ij}+l_{2ij}} \right) \right)^{-1}}} \quad (39)$$

Equivalent ν_{12}

$$\nu_{12eq} = -\frac{1}{L} \sum_{j=1}^n \frac{\sum_{i=1}^m (l_{ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\sum_{i=1}^m \frac{(\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})}{\cos \alpha_{ij} \cos \beta_{ij}}} \quad (40)$$

 Equivalent ν_{21}

$$\nu_{21eq} = -\frac{L}{\sum_{j=1}^n \frac{\sum_{i=1}^m (l_{ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\frac{l_{1ij}^2 l_{2ij}^2 (l_{ij} + l_{2ij}) \cos \alpha_{ij} \cos \beta_{ij} (\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})}{(l_{ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})^2 \left(l_{3ij}^2 \cos^2 \gamma_{ij} \left(l_{3ij} + \frac{l_{ij} l_{2ij}}{l_{ij} + l_{2ij}} \right) + \frac{l_{1ij}^2 l_{2ij}^2 (l_{ij} + l_{2ij}) \cos^2 \alpha_{ij} \cos^2 \beta_{ij}}{(l_{ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})^2} \right)}}}} \quad (41)$$

- According to the notations used for a regular lattice by Gibson and Ashby (1999), the notations for lattices without any structural irregularity can be expressed as: $L = n(h + l \sin \theta)$; $l_{1ij} = l_{2ij} = l_{3ij} = l$; $\alpha_{ij} = \theta$; $\beta_{ij} = 180^\circ - \theta$; $\gamma_{ij} = 90^\circ$, for all i and j .
- Using these transformations in case of the spatial variation of only material properties, the closed-form formulae for compound variation of material and geometric properties (equations 37–39) can be reduced to:

$$E_{1v} = \kappa_1 \left(\frac{t}{l}\right)^3 \frac{\cos \theta}{\left(\frac{h}{l} + \sin \theta\right) \sin^2 \theta} \quad (42)$$

$$E_{2v} = \kappa_2 \left(\frac{t}{l}\right)^3 \frac{\left(\frac{h}{l} + \sin \theta\right)}{\cos^3 \theta} \quad (43)$$

$$\text{and } G_{12v} = \kappa_2 \left(\frac{t}{l}\right)^3 \frac{\left(\frac{h}{l} + \sin \theta\right)}{\left(\frac{h}{l}\right)^2 \left(1 + 2\frac{h}{l}\right) \cos \theta} \quad (44)$$

- The multiplication factors κ_1 and κ_2 arising due to the consideration of spatially random variation of intrinsic material properties can be expressed as

$$\kappa_1 = \frac{m}{n} \sum_{j=1}^n \frac{1}{\sum_{i=1}^m \frac{1}{E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega} \right)}} \quad (45)$$

$$\text{and } \kappa_2 = \frac{n}{m} \frac{1}{\sum_{j=1}^n \frac{1}{\sum_{i=1}^m E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega} \right)}} \quad (46)$$

- In the special case when $\omega \rightarrow 0$ and there is no spatial variabilities in the material properties of the lattice, all viscoelastic material properties become identical (i.e. $E_{sij} = E_s$, $\mu_{ij} = \mu$ and $\epsilon_{ij} = \epsilon$ for $i = 1, 2, 3, \dots, m$ and $j = 1, 2, 3, \dots, n$) and subsequently the amplitude of κ_1 and κ_2 becomes exactly 1. This confirms that the expressions in 45 and 46 give the necessary generalisations of the classical expressions of Gibson and Ashby (1999) through 42–44.

- In case of only spatially random variation of structural geometry but constant viscoelastic material properties (i.e. $E_{sij} = E_S$, $\mu_{ij} = \mu$ and $\epsilon_{ij} = \epsilon$ for $i = 1, 2, 3, \dots, m$ and $j = 1, 2, 3, \dots, n$) the 37–39 lead to

$$E_{1v} = E_S \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \zeta_1 \quad (47)$$

$$E_{2v} = E_S \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \zeta_2 \quad (48)$$

$$G_{12v} = E_S \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \zeta_3 \quad (49)$$

- The random coefficients ζ_i ($i = 1, 2, 3$) are

$$\zeta_1 = \frac{t^3}{L} \sum_{j=1}^n \frac{\sum_{i=1}^m (l_{ij} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})}{\sum_{i=1}^m \frac{l_{1ij}^2 l_{2ij}^2 (l_{ij} + l_{2j}) (\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})^2}{(l_{ij} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})^2}} \quad (50)$$

$$\zeta_2 = \frac{Lt^3}{\sum_{j=1}^n \frac{\sum_{i=1}^m (l_{ij} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})}{\sum_{i=1}^m \left(l_{3ij}^2 \cos^2 \gamma_{ij} \left(l_{3ij} + \frac{l_{1ij} l_{2ij}}{l_{ij} + l_{2j}} \right) + \frac{l_{1ij}^2 l_{2ij}^2 (l_{ij} + l_{2j}) \cos^2 \alpha_{ij} \cos^2 \beta_{ij}}{(l_{ij} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})^2} \right)^{-1}}} \quad (51)$$

$$\zeta_3 = \frac{Lt^3}{\sum_{j=1}^n \frac{\sum_{i=1}^m (l_{ij} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})}{\sum_{i=1}^m \left(l_{3ij}^2 \sin^2 \gamma_{ij} \left(l_{3ij} + \frac{l_{1ij} l_{2ij}}{l_{ij} + l_{2j}} \right) \right)^{-1}}} \quad (52)$$

- The geometric notations for regular lattices can be expressed as:
 $L = n(h + l \sin \theta)$; $l_{1ij} = l_{2ij} = l_{3ij} = l$; $\alpha_{ij} = \theta$; $\beta_{ij} = 180^\circ - \theta$; $\gamma_{ij} = 90^\circ$, for all i and j . Using these transformations, the expressions of in-plane elastic moduli for regular hexagonal lattices (without the viscoelastic effect) can be obtained.
- The in-plane Young's moduli and shear modulus (viscosity dependent in-plane elastic properties) can be expressed as

$$E_{1v} = E_s \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \left(\frac{t}{l} \right)^3 \frac{\cos \theta}{\left(\frac{h}{l} + \sin \theta \right) \sin^2 \theta} \quad (53)$$

$$E_{2v} = E_s \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \left(\frac{t}{l} \right)^3 \frac{\left(\frac{h}{l} + \sin \theta \right)}{\cos^3 \theta} \quad (54)$$

$$G_{12v} = E_s \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \left(\frac{t}{l} \right)^3 \frac{\left(\frac{h}{l} + \sin \theta \right)}{\left(\frac{h}{l} \right)^2 \left(1 + 2 \frac{h}{l} \right) \cos \theta} \quad (55)$$

- The amplitude of the elastic moduli obtained based on the above expressions converge to the closed-form equation provided by [39] in the limiting case of $\omega \rightarrow 0$.

- In the case of regular uniform lattices with $\theta = 30^\circ$, we have

$$E_{1v} = E_{2v} = 2.3E_S \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \left(\frac{t}{l} \right)^3 \quad (56)$$

- Similarly, in the case of shear modulus for regular uniform lattices ($\theta = 30^\circ$)

$$G_{12v} = 0.57E_S \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \left(\frac{t}{l} \right)^3 \quad (57)$$

- Regular viscoelastic lattices satisfy the reciprocal theorem

$$E_{2v}\nu_{12v} = E_{1v}\nu_{21v} = E_S \left(1 + \epsilon \frac{i\omega}{\mu + i\omega} \right) \left(\frac{t}{l} \right)^3 \frac{1}{\sin \theta \cos \theta} \quad (58)$$

- Correlated structural and material attributes can be modelled random fields $\mathcal{H}(\mathbf{x}, \theta)$.
- The traditional way of dealing with random field is to discretise the random field into finite number of random variables. The available schemes for discretising random fields can be broadly divided into three groups: (1) point discretisation (e.g., midpoint method, shape function method, integration point method, optimal linear estimate method); (2) average discretisation method (e.g., spatial average, weighted integral method), and (3) series expansion method (e.g., orthogonal series expansion).
- An advantageous alternative for discretising $\mathcal{H}(\mathbf{x}, \theta)$ is to represent it in a generalised Fourier type of series as, often termed as Karhunen-Loève (KL) expansion.

- Suppose, $\mathcal{H}(\mathbf{x}, \theta)$ is a random field with covariance function $\Gamma_{\mathcal{H}}(\mathbf{x}_1, \mathbf{x}_2)$ defined in the probability space $(\Theta, \mathcal{F}, \mathcal{P})$. The KL expansion for $\mathcal{H}(\mathbf{x}, \theta)$ takes the following form

$$\mathcal{H}(\mathbf{x}, \theta) = \bar{\mathcal{H}}(\mathbf{x}) + \sum_{i=1}^{\infty} \sqrt{\lambda_i} \xi_i(\theta) \psi_i(\mathbf{x}) \quad (59)$$

where $\{\xi_i(\theta)\}$ is a set of uncorrelated random variables.

- $\{\lambda_i\}$ and $\{\psi_i(\mathbf{x})\}$ are the eigenvalues and eigenfunctions of the covariance kernel $\Gamma_{\mathcal{H}}(\mathbf{x}_1, \mathbf{x}_2)$, satisfying the integral equation

$$\int_{\mathbb{R}^N} \Gamma_{\mathcal{H}}(\mathbf{x}_1, \mathbf{x}_2) \psi_i(\mathbf{x}_1) d\mathbf{x}_1 = \lambda_i \psi_i(\mathbf{x}_2) \quad (60)$$

- In practise, the infinite series of 59 must be truncated, yielding a truncated KL approximation

$$\tilde{\mathcal{H}}(\mathbf{x}, \theta) \cong \bar{\mathcal{H}}(\mathbf{x}) + \sum_{i=1}^M \sqrt{\lambda_i} \xi_i(\theta) \psi_i(\mathbf{x}) \quad (61)$$

- Gaussian and lognormal random fields have been considered. The covariance function is represented as:

$$\Gamma_{\alpha z} = \sigma_{\alpha z}^2 e^{(-|y_1 - y_2|/b_y) + (-|z_1 - z_2|/b_z)} \quad (62)$$

where b_y and b_z are the correlation parameters at y and z directions (that corresponds to direction - 1 and direction - 2 respectively). These quantities control the rate at which the covariance decays.

- In a two dimensional physical space the eigensolutions of the covariance function are obtained by solving the integral equation analytically

$$\lambda_i \psi_i(y_2, z_2) = \int_{-a_1}^{a_1} \int_{-a_2}^{a_2} \Gamma(y_1, z_1; y_2, z_2) \psi_i(y_1, z_1) dy_1 dz_1 \quad (63)$$

where $-a_1 \leq y \leq a_1$ and $-a_2 \leq z \leq a_2$.

- Assume the eigen-solutions are separable in y and z directions, i.e.

$$\psi_i(y_2, z_2) = \psi_i^{(y)}(y_2) \psi_i^{(z)}(z_2) \quad (64)$$

$$\lambda_i(y_2, z_2) = \lambda_i^{(y)}(y_2) \lambda_i^{(z)}(z_2) \quad (65)$$

- The solution of the integral equation reduces to the product of the solutions of two equations of the form

$$\lambda_i^{(y)} \psi_i^{(y)}(y_1) = \int_{-a_1}^{a_1} e^{(-|y_1 - y_2|/b_y)} \psi_i^{(y)}(y_2) dy_2 \quad (66)$$

- The solution of this equation, which is the eigensolution (eigenvalues and eigenfunctions) of an exponential covariance kernel for a one-dimensional random field is obtained as

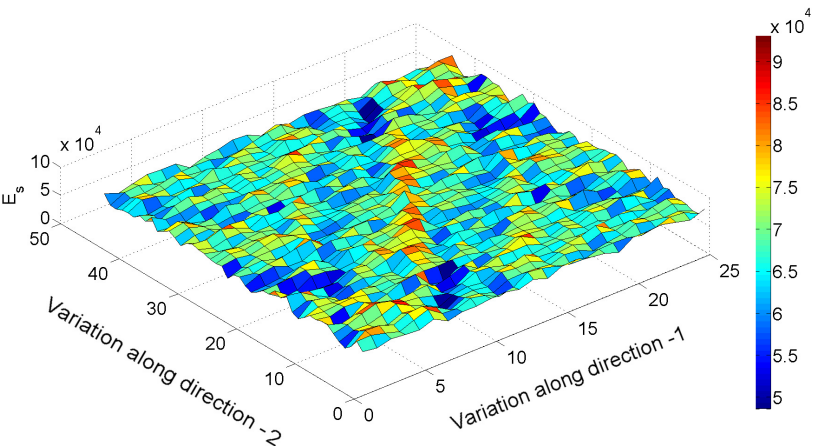
$$\begin{cases} \psi_i(\zeta) = \frac{\cos(\omega_i \zeta)}{\sqrt{a + \frac{\sin(2\omega_i a)}{2\omega_i}}} & \lambda_i = \frac{2\sigma_{\alpha z}^2 b}{\omega_i^2 + b^2} \quad \text{for } i \text{ odd} \\ \psi_i(\zeta) = \frac{\sin(\omega_i^* \zeta)}{\sqrt{a - \frac{\sin(2\omega_i^* a)}{2\omega_i^*}}} & \lambda_i^* = \frac{2\sigma_{\alpha z}^2 b}{\omega_i^{*2} + b^2} \quad \text{for } i \text{ even} \end{cases} \quad (67)$$

where $b = 1/b_y$ or $1/b_z$ and $a = a_1$ or a_2 . ζ can be either y or z and ω_i presents the period of the random field.

- The final eigenfunctions are given by

$$\psi_k(y, z) = \psi_i^{(y)}(y) \psi_i^{(z)}(z) \quad (68)$$

Samples of the random fields



Spatial variability of the intrinsic elastic modulus (E_s) with $\Delta_m = 0.002$

- To define the degree of irregularity, it is assumed that each connecting node of the lattice moves randomly within a certain radius (r_d) around the respective node corresponding to the regular deterministic configuration. For physically realistic variabilities, it is considered that a given node do not cross a neighbouring node, that is

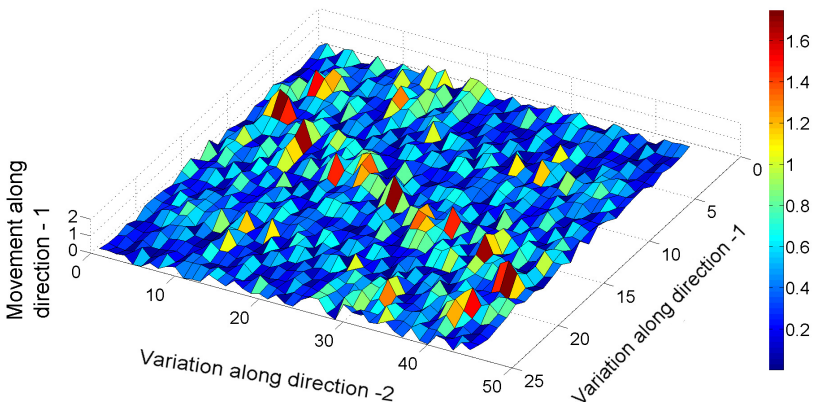
$$r_d < \min \left(\frac{h}{2}, \frac{l}{2}, l \cos \theta \right) \quad (69)$$

- In each realization of the Monte Carlo simulation, all the nodes of the lattice move simultaneously to new random locations within the specified circular bounds. Thus, the degree of irregularity (r) is defined as a non-dimensional ratio of the area of the circle and the area of one regular hexagonal unit as

$$r = \frac{\pi r_d^2 \times 100}{2l \cos \theta (h + l \sin \theta)} \quad (70)$$

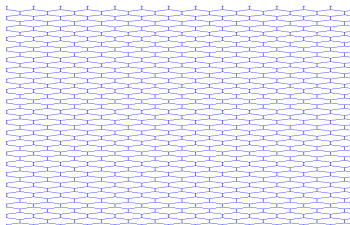
- The degree of irregularity (r) has been expressed as percentage values for presenting the results.

Samples of the random fields

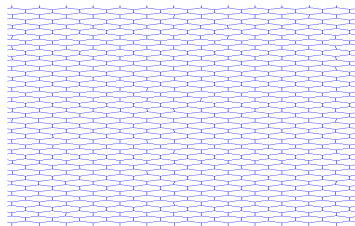


Movement of the top vertices of a tessellating hexagonal unit cell with respect to the corresponding deterministic locations ($r = 6$)

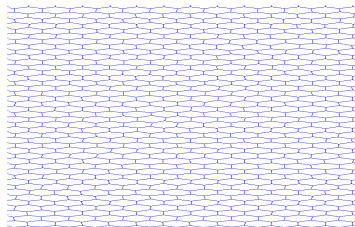
Random geometric configurations



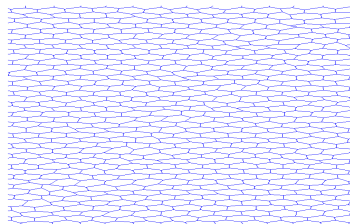
(a)




(b)



(c)



(d)

Structural configurations for a single random realisation of an irregular hexagonal lattice considering deterministic cell angle $\theta = 30^\circ$ and $h/l = 1$: (a)  $r = 0$ (b) $r = 2$ (c) $r = 4$ (d) $r = 6$

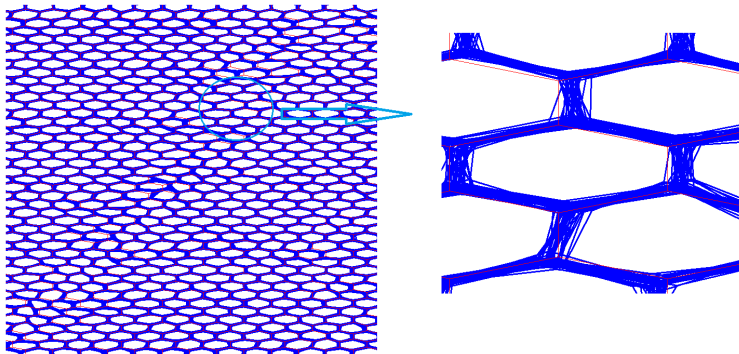


Figure: Simulation bound of the structural configuration of an irregular hexagonal lattice for multiple random realisations considering $\theta = 30^\circ$, $h/l = 1$ and $r = 6$. The regular configuration is presented using red colour.

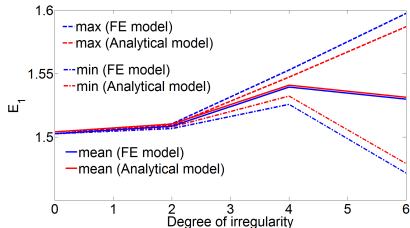
- In randomly inhomogeneous correlated system, spatial variability of the stochastic structural attributes are accounted, wherein each sample of the Monte Carlo simulation includes the spatially random distribution of structural and materials attributes with a rule of correlation.
- The spatial variability in structural and material properties (E_s , μ and ϵ) are physically attributed by **degree of structural irregularity (r)** and **degree of material property variation (Δ_m)** respectively.
- As the two Young's moduli and shear modulus for low density lattices are proportional to $E_s\rho^3$ [40], the **non-dimensional results** for in-plane elastic moduli E_1 , E_2 , and G_{12} , unless otherwise mentioned, are presented as:

$$\bar{E}_1 = \frac{E_{1eq}}{E_s\rho^3}, \quad \bar{E}_2 = \frac{E_{2eq}}{E_s\rho^3}$$

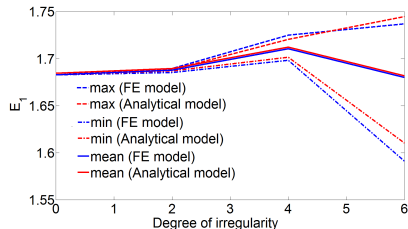
$$\bar{G}_{12} = \frac{G_{12eq}}{E_s\rho^3}$$

- ρ is the relative density of the lattice (defined as a ratio of the planar area of solid to the total planar area of the lattice).

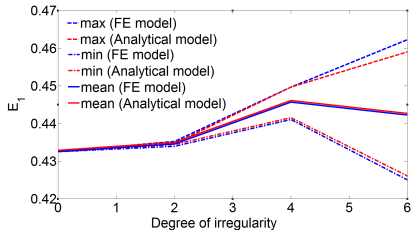
Spatially correlated irregular elastic lattices: E_1



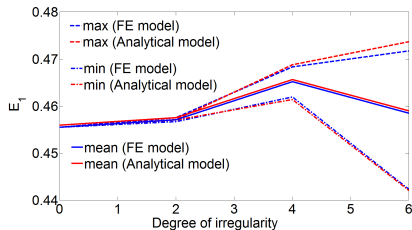
(a) $\theta = 30^\circ$; $\frac{h}{l} = 1$



(b) $\theta = 30^\circ$; $\frac{h}{l} = 1.5$



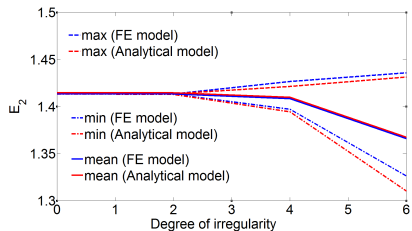
(c) $\theta = 45^\circ$; $\frac{h}{l} = 1$



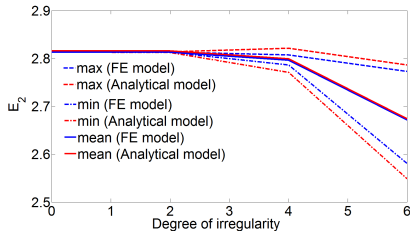
(d) $\theta = 45^\circ$; $\frac{h}{l} = 1.5$

Figure: Effective Young's modulus (E_1) of irregular lattices

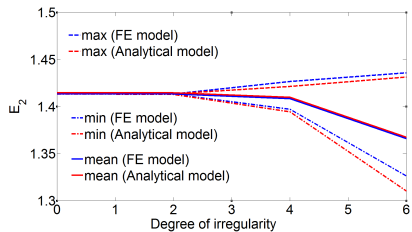
Spatially correlated irregular elastic lattices: E_2



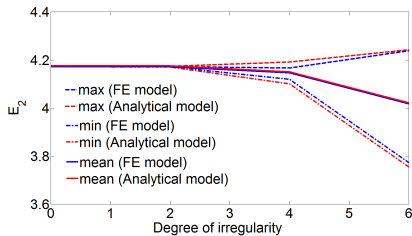
(a) $\theta = 30^\circ$; $\frac{h}{l} = 1$



(b) $\theta = 30^\circ$; $\frac{h}{l} = 1.5$

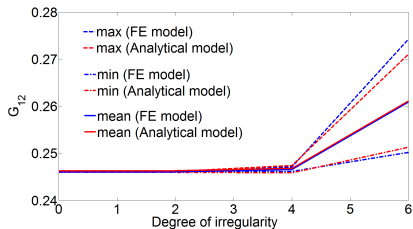


(c) $\theta = 45^\circ$; $\frac{h}{l} = 1$

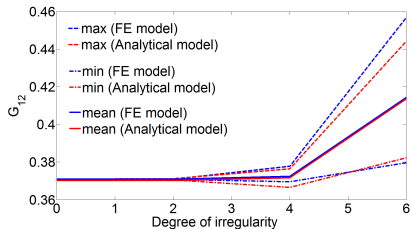


(d) $\theta = 45^\circ$; $\frac{h}{l} = 1.5$

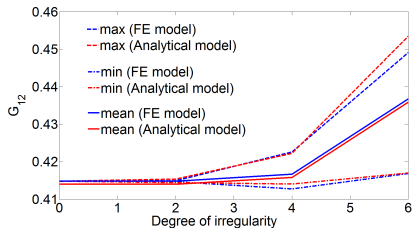
Spatially correlated irregular elastic lattices: G_{12}



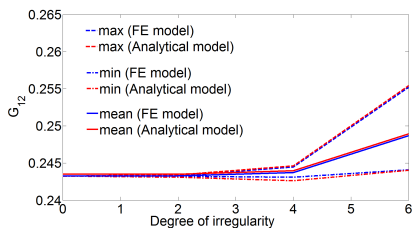
(a) $\theta = 30^\circ$; $\frac{h}{T} = 1$



(b) $\theta = 30^\circ$; $\frac{h}{T} = 1.5$

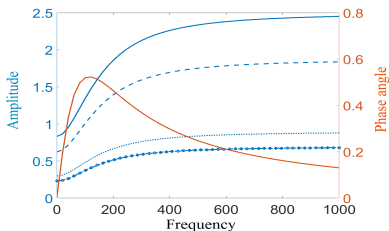


(c) $\theta = 45^\circ$; $\frac{h}{T} = 1$

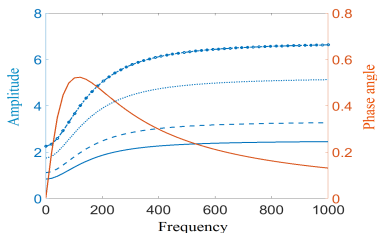


(d) $\theta = 45^\circ$; $\frac{h}{T} = 1.5$

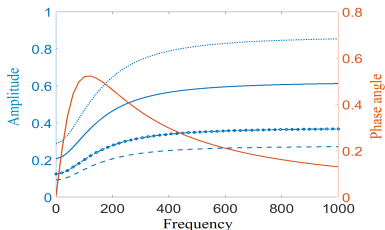
Viscoelastic properties of regular lattices: E_1, E_2, G_{12}



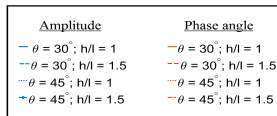
(a)



(b)

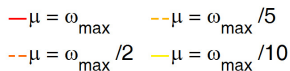
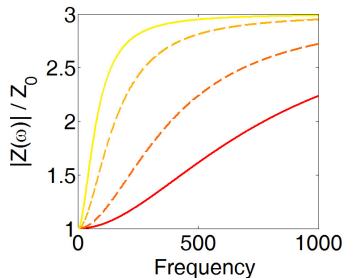


(c)

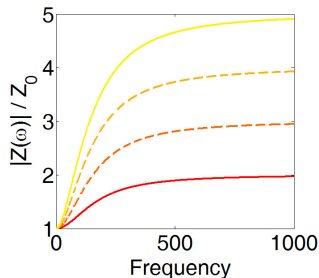


(a) Effect of viscoelasticity on the magnitude and phase angle of E_1 for regular hexagonal lattices (b) Effect of viscoelasticity on the magnitude and phase angle of E_2 for regular hexagonal lattices (c) Effect of viscoelasticity on the magnitude and phase angle of G_{12} for regular hexagonal lattices

Viscoelastic properties of regular lattices



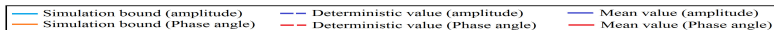
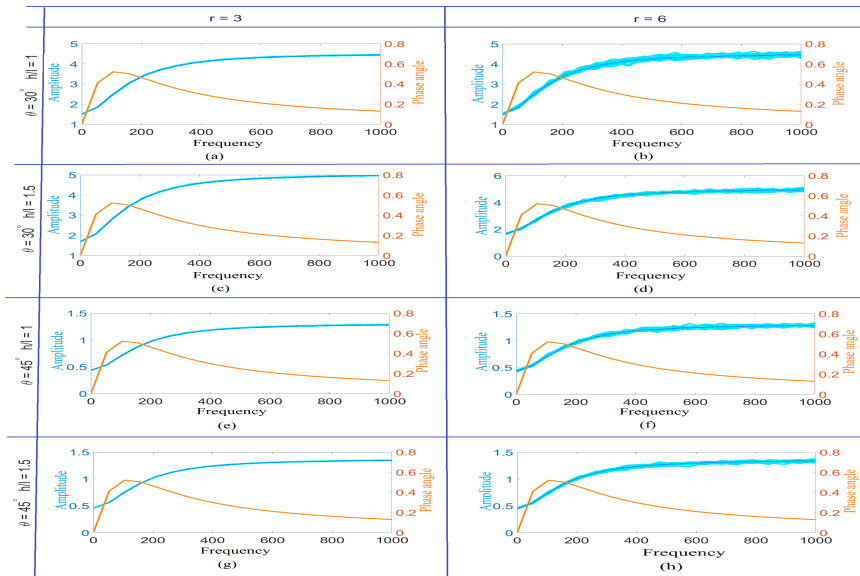
(a)



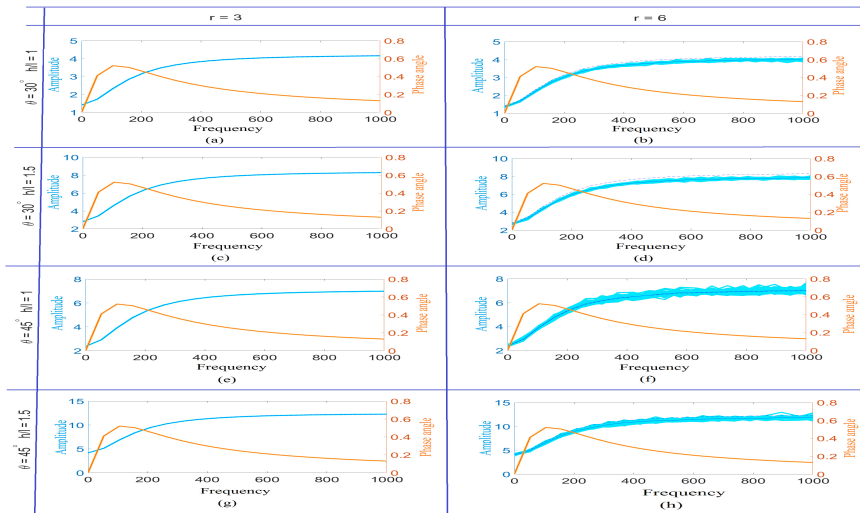
(b)

(a) Effect of variation of μ on the viscoelastic modulus of regular hexagonal lattices (considering a constant value of $\epsilon = 2$) (b) Effect of variation of ϵ on the viscoelastic modulus of regular hexagonal lattices (considering a constant value of $\mu = \omega_{\max} / 5$). Here Z represents the viscoelastic moduli (i.e. E_1 , E_2 and G_{12}) and Z_0 is the corresponding elastic modulus value for $\omega = 0$.

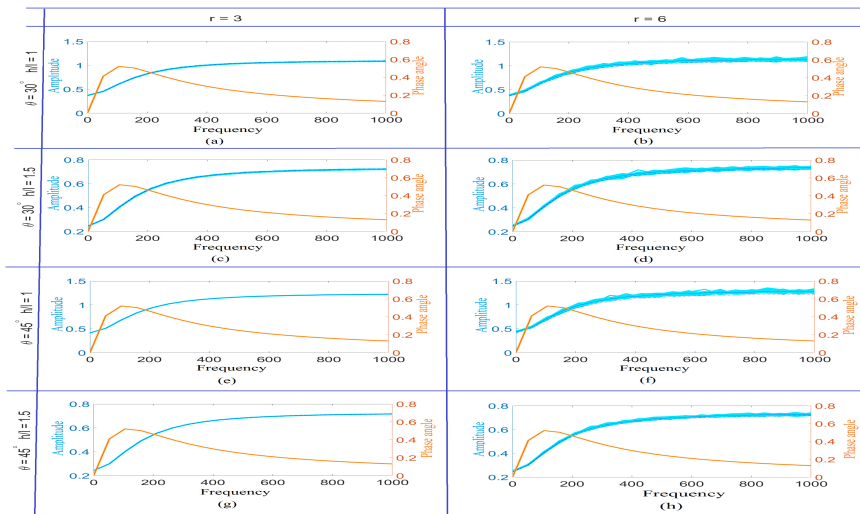
Spatially correlated irregular viscoelastic lattices: E_1



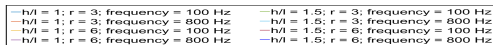
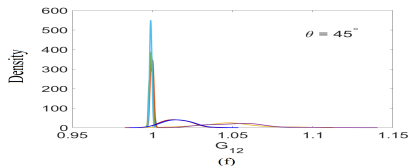
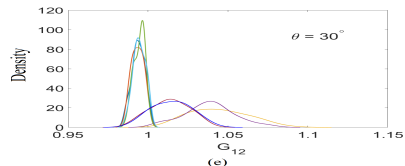
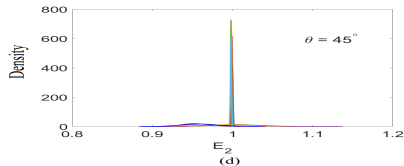
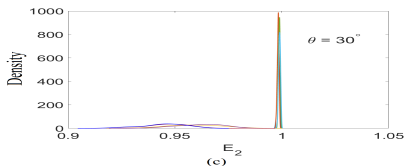
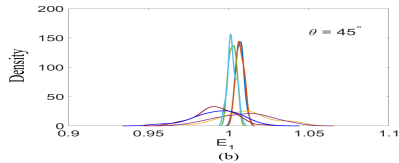
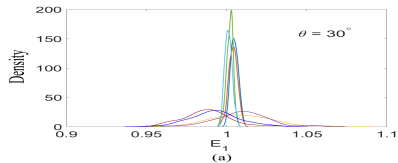
Spatially correlated irregular viscoelastic lattices: E_2



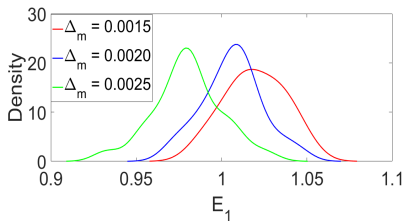
Spatially correlated irregular elastic lattices: G_{12}



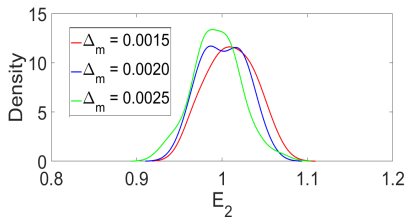
Probability density function: random geometry



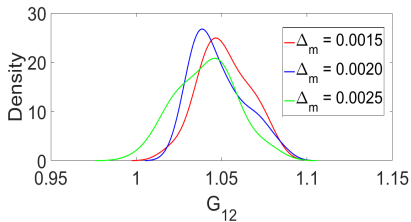
Probability density function: random material property



(a)



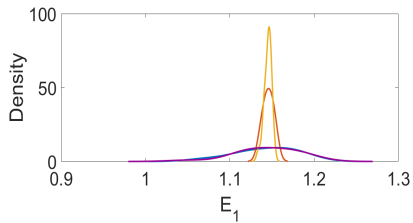
(b)



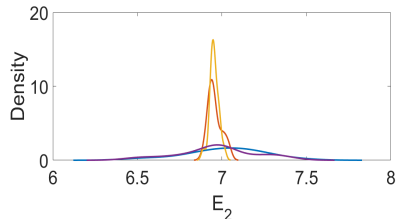
(c)

Probability density function plots for the amplitude of the elastic moduli considering randomly inhomogeneous form of stochasticity for different values of Δ_m (i.e. coefficient of variation for spatially random correlated material properties, such as E_S , μ and ϵ). Results are presented as a ratio of the values corresponding to irregular configurations and respective deterministic values (for a frequency of 800 Hz).

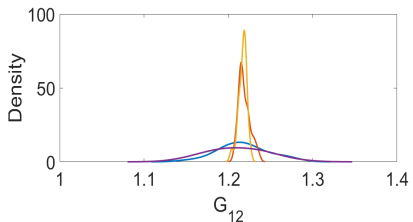
Combined material and geometric uncertainty



(a)



(b)



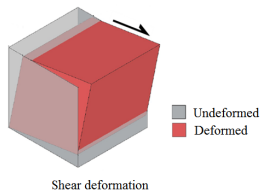
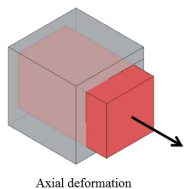
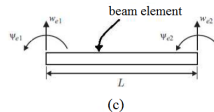
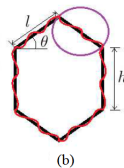
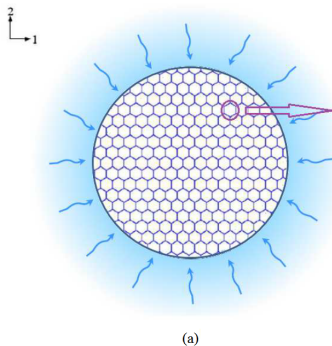
(c)

- Stochasticity in structural attributes
- Stochasticity in intrinsic elastic modulus (E_s)
- Stochasticity in viscoelastic parameters (μ and ϵ)
- Combined stochasticity

Probabilistic descriptions for the amplitudes of three effective viscoelastic properties corresponding to a frequency of 800 Hz considering individual and compound effect of stochasticity in material and structural attributes with $\Delta_{COV} = 0.006$



Cellular structure in a dynamic environment



(d)

- The equation of motion of free vibration of a damped beam can be expressed as

$$EI \frac{\partial^4 V(x, t)}{\partial x^4} + \hat{c}_1 \frac{\partial^5 V(x, t)}{\partial x^4 \partial t} + m \frac{\partial^2 V(x, t)}{\partial t^2} + \hat{c}_2 \frac{\partial V(x, t)}{\partial t} = 0 \quad (71)$$

- From the equation of motion, the dynamic stiffness matrix of a single beam element can be derived in a closed form as:

$$\mathbf{D}(\omega) = \frac{EIb}{(cC - 1)} \begin{bmatrix} -b^2 (cS + Cs) & -sbS & b^2 (S + s) & -b(C - c) \\ -sbS & -Cs + cS & b(C - c) & -S + s \\ b^2 (S + s) & b(C - c) & -b^2 (cS + Cs) & sbS \\ -b(C - c) & -S + s & sbS & -Cs + cS \end{bmatrix} \quad (72)$$

where

$$C = \cosh(bL), c = \cos(bL), S = \sinh(bL), s = \sin(bL), b^4 = \frac{m\omega^2 (1 - i\zeta_m/\omega)}{EI(1 + i\omega\zeta_k)} \quad (73)$$

The quantities ζ_k and ζ_m are stiffness and mass proportional damping factors.



- The **mass distribution** of the element is treated in an exact manner in deriving the element dynamic stiffness matrix.
- The dynamic stiffness matrix of one-dimensional structural elements, taking into account the effects of flexure, torsion, axial and shear deformation, and damping, is **exactly determinable**, which, in turn, enables the exact vibration analysis by an inversion of the global dynamic stiffness matrix.
- The method **does not employ eigenfunction expansions** and, consequently, a major step of the traditional finite element analysis, namely, the determination of natural frequencies and mode shapes, is eliminated which automatically avoids the errors due to series truncation.
- The **damping** within the system can be incorporated in a rigorous manner using complex algebra.
- The method is essentially a **frequency-domain approach** suitable for steady state harmonic or stationary random excitation problems.
- The static stiffness matrix and the consistent mass matrix appear as the **first two terms in the Taylor expansion** of the dynamic stiffness matrix in the frequency parameter.

Frequency-dependent elastic moduli

- The equivalent elastic moduli of the hexagonal cellular materials (Eqs. 1-5) were obtained considering only the static deformation of a unit cell.
- Exploiting the physical interpretation of the elements of the dynamic stiffness matrix,
- the analytical expressions of the frequency dependent in-plane elastic moduli of the hexagonal lattice are obtained.
- Based on the concept of a unit cell based approach, one hexagonal unit is analysed to derive the expressions of in-plane elastic moduli of the lattice material:

Equivalent $E_1(\omega)$

$$\begin{aligned} E_1(\omega) &= \frac{D_{33} l \cos \theta}{(h + l \sin \theta) \bar{b} \sin^2 \theta} \\ &= \frac{Et^3 l \cos \theta b^3 (\cos(bl) \sinh(bl) + \cosh(bl) \sin(bl))}{12(h + l \sin \theta) \sin^2 \theta (1 - \cos(bl) \cosh(bl))} \end{aligned} \quad (74)$$

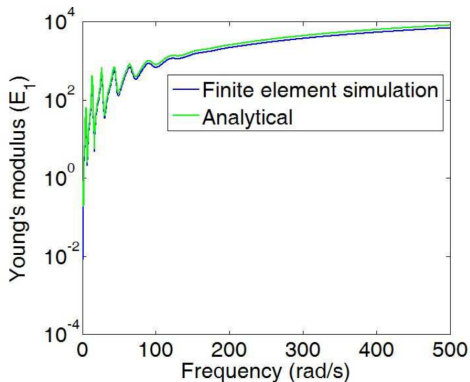
Equivalent $E_2(\omega)$

$$\begin{aligned}
 E_2(\omega) &= \frac{D_{33}(h + l \sin \theta)}{l\bar{b} \cos^3 \theta} \\
 &= \frac{Et^3(h + l \sin \theta)b^3 (\cos(bl) \sinh(bl) + \cosh(bl) \sin(bl))}{12l \cos^3 \theta (1 - \cos(bl) \cosh(bl))}
 \end{aligned} \tag{75}$$

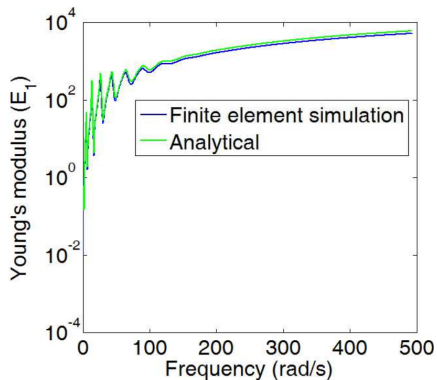
Equivalent $G_{12}(\omega)$

$$\begin{aligned}
 G_{12}(\omega) &= \frac{(h + l \sin \theta)}{2l\bar{b} \cos \theta} \frac{1}{\left(-\frac{h^2}{4lD_{43}^s} + \frac{2}{\left(D_{33}^v - \frac{(D_{34}^v)^2}{D_{44}^v} \right)} \right)} \\
 &= \frac{Et^3(h + l \sin \theta)b^3 \sin(bl) \sinh(bl) (1 + \cos(bh/2) \cosh(bh/2))}{6l \cos \theta [h^2b(1 - \cos(bl) \cosh(bl)) (1 + \cos(bh/2) \cosh(bh/2)) \\
 &\quad + 8l \sin(bl) \sinh(bl) (\cosh(bh/2) \sin(bh/2) - \sinh(bh/2) \cos(bh/2))] }
 \end{aligned} \tag{76}$$

Frequency dependent Young's modulus (E_1)



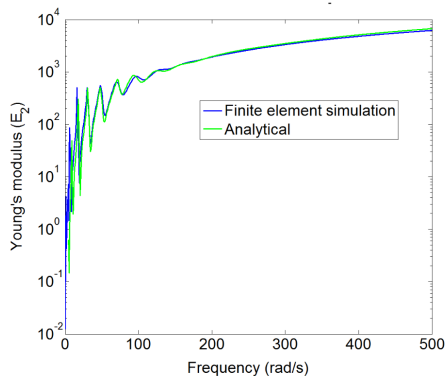
(a) $h/l = 1$



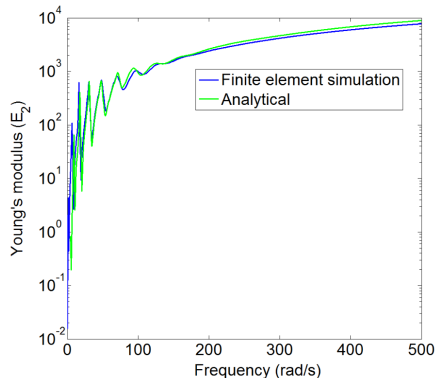
(b) $h/l = 1.5$

Frequency dependent Young's modulus (E_1) of regular hexagonal lattices with $\theta = 30^\circ$ and $\zeta_m = 0.05$ and $\zeta_k = 0.002$.

Frequency dependent Young's modulus (E_2)



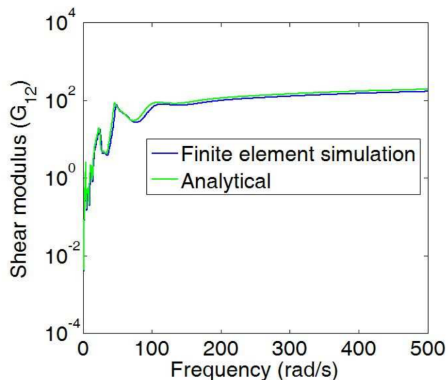
(a) $h/l = 1$



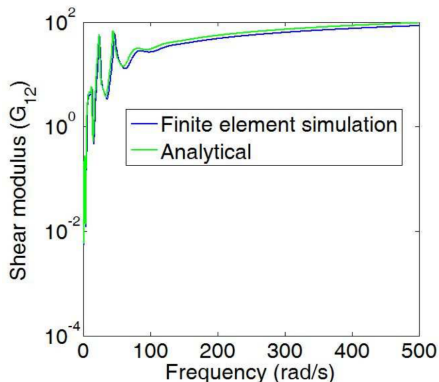
(b) $h/l = 1.5$

Frequency dependent Young's modulus (E_2) of regular hexagonal lattices with $\theta = 30^\circ$ and $\zeta_m = 0.05$ and $\zeta_k = 0.002$.

Frequency dependent Young's modulus (G_{12})



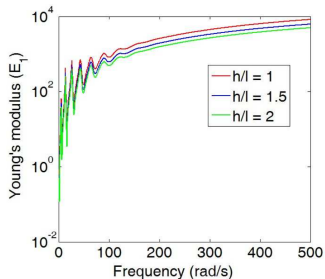
(a) $h/l = 1$



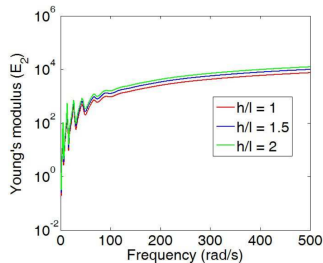
(b) $h/l = 1.5$

Frequency dependent shear modulus (G_{12}) of regular hexagonal lattices with $\theta = 30^\circ$ and $\zeta_m = 0.05$ and $\zeta_k = 0.002$.

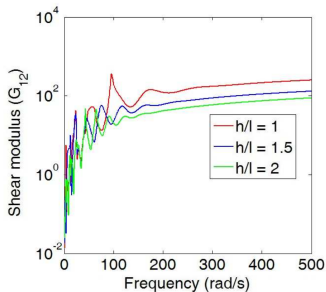
Variation of the Elastic moduli with different h/l ratios



(a) Young's modulus E_1

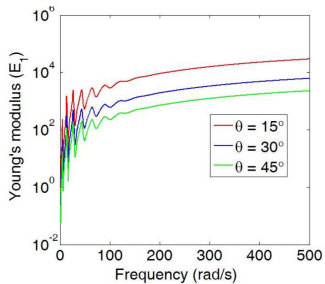


(b) Young's modulus E_2

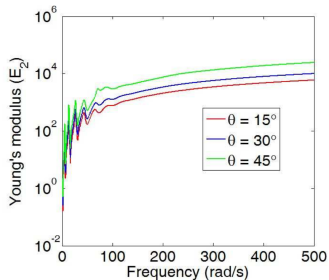


(c) Shear modulus G_{12}

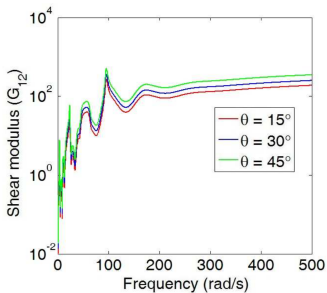
Variation of the Elastic moduli with different with cell angle θ



(a) Young's modulus E_1

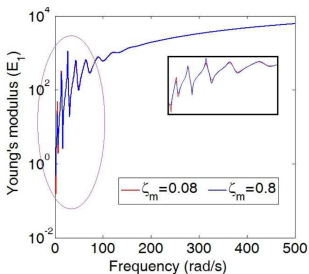


(b) Young's modulus E_2

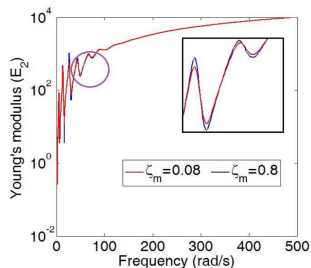


(c) Shear modulus G_{12}

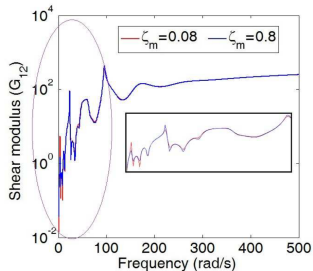
Variation of the Elastic moduli with different mass proportional damping factors



(a) Young's modulus E_1

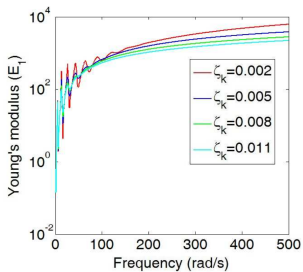


(b) Young's modulus E_2

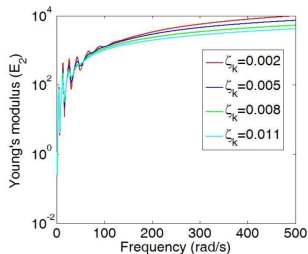


(c) Shear modulus G_{12}

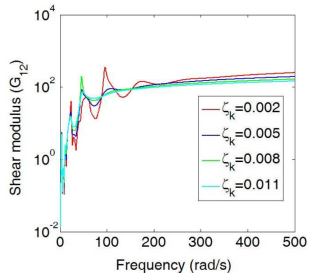
Variation of the Elastic moduli with different stiffness proportional damping factors



(a) Young's modulus E_1



(b) Young's modulus E_2



(c) Shear modulus G_{12}

- Considering the static case, that is, when the frequency $\omega \rightarrow 0$, we have

$$\lim_{\omega \rightarrow 0} D_{33} = 12 \frac{EI}{\beta^3} = 12 \left(\frac{1}{12} E \bar{b} t^3 \right) \frac{1}{\beta^3} = E \bar{b} \left(\frac{t}{l} \right)^3 \quad (77)$$

- Substituting this in the expressions of $E_1(\omega)$ and $E_2(\omega)$ in equations (92) and (93) we have

$$\begin{aligned} \lim_{\omega \rightarrow 0} E_1(\omega) &= \frac{l \cos \theta}{(h + l \sin \theta) \bar{b} \sin^2 \theta} \lim_{\omega \rightarrow 0} D_{33} \\ &= E \left(\frac{t}{l} \right)^3 \frac{l \cos \theta}{(h + l \sin \theta) \sin^2 \theta} \end{aligned} \quad (78)$$

$$\text{and } \lim_{\omega \rightarrow 0} E_2(\omega) = \frac{(h + l \sin \theta)}{l \bar{b} \cos^3 \theta} \lim_{\omega \rightarrow 0} D_{33} = E \left(\frac{t}{l} \right)^3 \frac{(h + l \sin \theta)}{l \cos^3 \theta} \quad (79)$$

- The above expressions match exactly with the original classical expressions of E_1 and E_2 derived by [39].

The static limit of the elastic moduli

- The shear modulus $G_{12}(\omega)$ given in (94) is a function of four dynamic stiffness coefficients. In the limiting case they become

$$\lim_{\omega \rightarrow 0} D_{43}^s = -6 \frac{EI}{l^2}, \quad \lim_{\omega \rightarrow 0} D_{33}^v = 96 \frac{EI}{h^3}, \quad \lim_{\omega \rightarrow 0} D_{34}^v = -24 \frac{EI}{h^2}, \quad \lim_{\omega \rightarrow 0} D_{44}^v = 8 \frac{EI}{h} \quad (80)$$

- Substituting these in (94) we have

$$\begin{aligned} \lim_{\omega \rightarrow 0} G_{12}(\omega) &= \frac{(h + l \sin \theta)}{2l\bar{b} \cos \theta} \lim_{\omega \rightarrow 0} \frac{1}{\left(-\frac{h^2}{4lD_{43}^s} + \frac{2}{\left(D_{33}^v - \frac{(D_{34}^v)^2}{D_{44}^v} \right)} \right)} \\ &= \frac{(h + l \sin \theta)}{2l\bar{b} \cos \theta} \frac{1}{(1/24) \frac{h^2 l}{EI} + (1/12) \frac{h^3}{EI}} = E \left(\frac{t}{l} \right)^3 \frac{\left(\frac{h}{l} + \sin \theta \right)}{\left(\frac{h}{l} \right)^2 (1 + 2 \frac{h}{l}) \cos \theta} \end{aligned} \quad (81)$$

- This shows that the shear modulus $G_{12}(\omega)$ also reduces to the classical equation in the static limit.
- These expressions should be viewed as the dynamic generalisation of the conventional equivalent elastic moduli of the hexagonal cellular material.

The undamped limit and the negative moduli

- The expressions of E_1 and E_2 are proportional to D_{33} , which is a complex frequency-dependent coefficient. Therefore, we study its behaviour in the undamped limit to understand the real part of E_1 and E_2 .
- Assuming no damping in the system, the parameter b in equation (73) becomes

$$b^4 = \frac{m\omega^2}{EI} \quad (82)$$

- Substituting this in the expression of D_{33} and expanding the expression by a Taylor series in the frequency parameter ω we have

$$D_{33} = 12 \frac{EI}{l^3} - \frac{13}{35} ml\omega^2 - \frac{59}{161700} \frac{l^5 m^2 \omega^4}{EI} - \frac{551}{794593800} \frac{l^9 m^3 \omega^6}{EI^2} + \dots \quad (83)$$

- Note that coefficients of some higher order terms of ω are negative. We observe that D_{33} appears as a multiplicative term in the expressions of $E_1(\omega)$ and $E_2(\omega)$ in equations (92) and (93) and the other terms are positive.
- Therefore, near the vicinity of $\omega \approx 0$, there exist some frequency beyond which the effective elastic moduli of honeycomb will be negative.

The undamped limit and the negative moduli

- Retaining up to terms of order ω^4 in equation (84), the critical value of ω can be obtained by setting $D_{33} = 0$ as

Negative E_1, E_2

$$D_{33} \approx 12 \frac{EI}{\beta^3} - \frac{13}{35} m l \omega^2 - \frac{59}{161700} \frac{l^5 m^2 \omega^4}{EI} = 0 \quad (84)$$

or $\omega_{E_1, E_2}^* = 5.598 \frac{1}{l^2} \sqrt{\frac{EI}{m}}$

- For lightly damped systems, beyond this frequency value, the equivalent Young's moduli E_1 and E_2 will be negative.
- Since the discovery of the Young's modulus over three centuries ago, it has been generally recognised as a positive quantity. When a dynamic equilibrium is considered, our results show that for cellular metamaterials the Young's moduli can be negative, contradicting notions established for centuries.
- Similar observation has been made in the context of acoustics metamaterials with sub-wavelength scale oscillators.

The undamped limit and the negative moduli

- For the shear modulus, it is also possible to expand the frequency dependent expression (94) to expand in a Taylor series in ω about $\omega = 0$ as

$$G_{12}(\omega) = \frac{(h + l \sin \theta)}{2l\bar{b} \cos \theta} \left[24 \frac{EI}{h^2 (2h + l)} - \frac{11}{420} \frac{m (9h^5 + 8l^5) \omega^2}{h^2 (2h + l)^2} \right. \\ \left. - \frac{1}{46569600} \frac{m^2 (55461h^9l - 191664h^5l^5 + 198912l^9h + 3111h^{10} + 14272l^{10}) \omega^4}{Elh^2 (2h + l)^3} \right] \dots \quad (85)$$

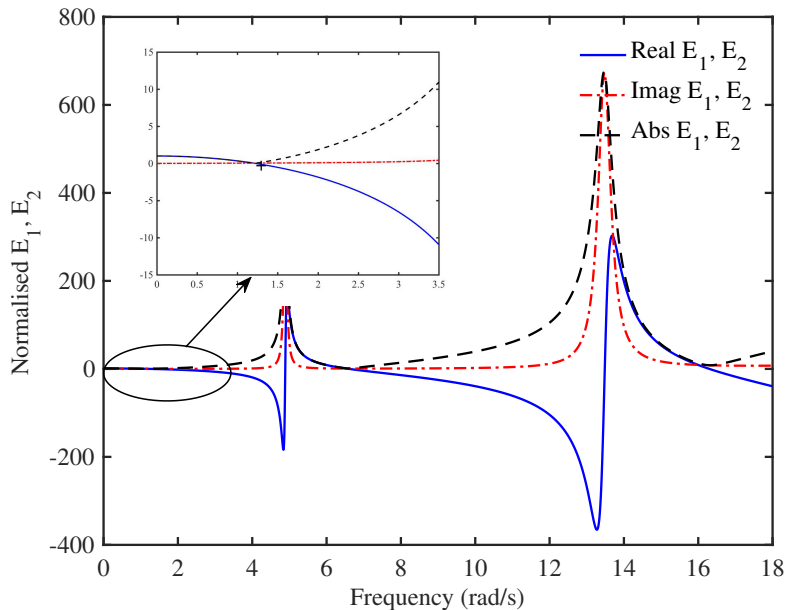
- Considering only up to the second-order terms, we obtain the following fundamental inequality regarding the frequency for negative value of G_{12}

Negative G_{12}

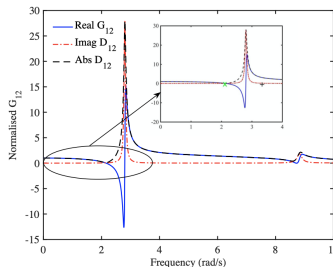
$$\frac{120}{\sqrt{160 + 75(h/l)^4}} \frac{1}{l^2} \sqrt{\frac{EI}{m}} < \omega_{G_{12}}^* < 30.2715 \sqrt{\frac{1 + 2(h/l)}{8 + 9(h/l)^5}} \frac{1}{l^2} \sqrt{\frac{EI}{m}} \quad (86)$$

- Unlike the equivalent expression for the Young's moduli E_1 and E_2 in (84), for the minimum frequency above which G_{12} becomes negative depends on the h/l ratio.

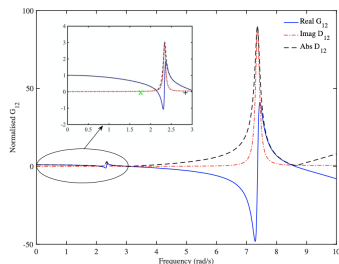
Negative E_1 and E_2



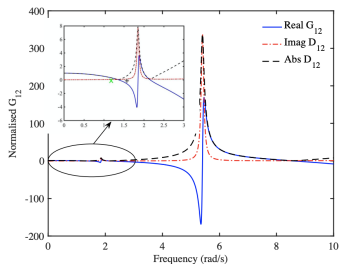
The real and imaginary parts and the amplitude of the normalised value of E_1 and E_2 as a function of frequency



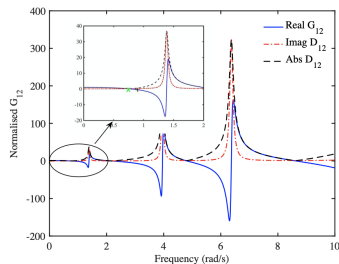
(a) $h/l = 0.5$, $2.0430 < \omega_{G_{12}}^* < 3.2502$



(b) $h/l = 1.0$, $1.7103 < \omega_{G_{12}}^* < 2.7783$



(c) $h/l = 1.5$, $1.1286 < \omega_{G_{12}}^* < 1.5139$



(d) $h/l = 2.0$, $0.71090 < \omega_{G_{12}}^* < 0.8596$

The real and imaginary parts and the amplitude of the normalised value of G_{12} as a function of frequency for four different values of h/l .

- The effect of viscoelasticity on irregular hexagonal lattices is investigated in frequency domain considering two different forms of irregularity in structural and material parameters (spatially uncorrelated and correlated).
- Spatially correlated structural and material attributes are considered to account for the effect of randomly inhomogeneous form of irregularity based on Karhunen-Loève expansion.
- The two Young's moduli and shear modulus are dependent on the viscoelastic parameters. Two in-plane Poisson's ratios depend only on structural geometry of the lattice structure.
- Using the principle of basic structural mechanics on a newly defined unit cell with a homogenisation technique, closed-form expressions have been obtained for E_1 , E_2 , ν_{12} , ν_{21} and G_{12} .
- Frequency dependent equivalent homogeneous properties were derived using the dynamic stiffness approach.
- The classical closed-form expressions for equivalent in-plane elastic properties of regular hexagonal lattice structures have been generalised to consider geometric and material irregularity, viscoelasticity and inertial effect.
- The new results reduce to classical formulae of Gibson and Ashby for the special case of no irregularities, no viscoelastic effect and the static limit.
- Beyond some frequency values derived here, the effective elastic moduli of the lattice can be negative.

Closed-form expressions: Elastic Moduli of disordered viscoelastic lattice

$$E_{1V}(\omega) = \frac{t^3}{L} \sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\sum_{i=1}^m \frac{l_{1ij}^2 l_{2ij}^2 (l_{1ij} + l_{2ij}) (\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})^2}{E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega}\right) ((l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})^2)}} \quad (87)$$

$$E_{2V}(\omega) = \frac{Lt^3}{\sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\sum_{i=1}^m E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega}\right) \left(l_{3ij}^2 \cos^2 \gamma_{ij} \left(l_{3ij} + \frac{l_{1ij} l_{2ij}}{l_{1ij} + l_{2ij}} \right) + \frac{l_{1ij}^2 l_{2ij}^2 (l_{1ij} + l_{2ij}) \cos^2 \alpha_{ij} \cos^2 \beta_{ij}}{(l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})^2} \right)^{-1}}} \quad (88)$$

$$G_{12V}(\omega) = \frac{Lt^3}{\sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1ij} \cos \alpha_{ij} - l_{2ij} \cos \beta_{ij})}{\sum_{i=1}^m E_{sij} \left(1 + \epsilon_{ij} \frac{i\omega}{\mu_{ij} + i\omega}\right) \left(l_{3ij}^2 \sin^2 \gamma_{ij} \left(l_{3ij} + \frac{l_{1ij} l_{2ij}}{l_{1ij} + l_{2ij}} \right) \right)^{-1}}} \quad (89)$$

Closed-form expressions: Poisson's ratios of disordered viscoelastic lattice

$$\nu_{12eq} = -\frac{1}{L} \sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1j} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})}{\sum_{i=1}^m \frac{(\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})}{\cos \alpha_{ij} \cos \beta_{ij}}} \quad (90)$$

$$\nu_{21eq} = -\frac{L}{\sum_{j=1}^n \frac{\sum_{i=1}^m (l_{1j} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})}{\frac{l_{1j}^2 l_{2j}^2 (l_{1j} + l_{2j}) \cos \alpha_{ij} \cos \beta_{ij} (\cos \alpha_{ij} \sin \beta_{ij} - \sin \alpha_{ij} \cos \beta_{ij})}{(l_{1j} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})^2 \left(l_{3j}^2 \cos^2 \gamma_{ij} \left(l_{3j} + \frac{l_{1j} l_{2j}}{l_{1j} + l_{2j}} \right) + \frac{l_{1j}^2 l_{2j}^2 (l_{1j} + l_{2j}) \cos^2 \alpha_{ij} \cos^2 \beta_{ij}}{(l_{1j} \cos \alpha_{ij} - l_{2j} \cos \beta_{ij})^2} \right)}}}} \quad (91)$$

$$\begin{aligned}
 E_1(\omega) &= \frac{D_{33} l \cos \theta}{(h + l \sin \theta) \bar{b} \sin^2 \theta} \\
 &= \frac{Et^3 l \cos \theta b^3 (\cos(bl) \sinh(bl) + \cosh(bl) \sin(bl))}{12(h + l \sin \theta) \sin^2 \theta (1 - \cos(bl) \cosh(bl))}
 \end{aligned} \tag{92}$$

$$\begin{aligned}
 E_2(\omega) &= \frac{D_{33}(h + l \sin \theta)}{l \bar{b} \cos^3 \theta} \\
 &= \frac{Et^3 (h + l \sin \theta) b^3 (\cos(bl) \sinh(bl) + \cosh(bl) \sin(bl))}{12l \cos^3 \theta (1 - \cos(bl) \cosh(bl))}
 \end{aligned} \tag{93}$$

$$\begin{aligned}
 G_{12}(\omega) &= \frac{(h + l \sin \theta)}{2l \bar{b} \cos \theta} \frac{1}{\left(-\frac{h^2}{4lD_{43}^s} + \frac{2}{\left(D_{33}^v - \frac{(D_{34}^v)^2}{D_{44}^v} \right)} \right)} \\
 &= \frac{Et^3 (h + l \sin \theta) b^3 \sin(bl) \sinh(bl) (1 + \cos(bh/2) \cosh(bh/2))}{6l \cos \theta [h^2 b (1 - \cos(bl) \cosh(bl)) (1 + \cos(bh/2) \cosh(bh/2)) \\
 &\quad + 8l \sin(bl) \sinh(bl) (\cosh(bh/2) \sin(bh/2) - \sinh(bh/2) \cos(bh/2))] }
 \end{aligned} \tag{94}$$

Some of our papers on this topic

- 1 Mukhopadhyay, T., Adhikari, S., and Alu, A., "Theoretical limits for negative elastic moduli in subacoustic lattice materials", [Physical Review B](#), Vol. 99, 2019, pp. 094108.
- 2 Mukhopadhyay, T., Adhikari, S., and Alu, A., "Probing the frequency-dependent elastic moduli of lattice materials", [Acta Materialia](#), Vol. 165, No. 2, 2019, pp. 654-665.
- 3 Mukhopadhyay, T., Adhikari, S., and Batou, A., "Viscoelastic mechanical properties of irregular quasi-periodic lattices with spatially correlated material and structural attributes", [International Journal of Mechanical Science](#), Vol. 150, No. 1, 2019, pp. 784-806.
- 4 Mukhopadhyay, T., Mahata, T., Adhikari, S., and Zaeem, M. A., "Probing the shear modulus of two-dimensional multiplanar nanostructures and heterostructures", [Nanoscale](#), Vol. 10, No. 11, 2018, pp. 5280-5294.
- 5 Mukhopadhyay, T., Mahata, A., Adhikari, S. and Asle Zaeem, M., "Effective mechanical properties of multilayer nano-heterostructures", [Nature Scientific Reports](#), (2017), pp. 15818:1-13.
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- 8 Mukhopadhyay, T. and Adhikari, S., "Stochastic mechanics of metamaterials", [Composite Structures](#), 162[2] (2017), pp. 85-97.
- 9 Mukhopadhyay, T. and Adhikari, S., "Free vibration of sandwich panels with randomly irregular honeycomb core", [ASCE Journal of Engineering Mechanics](#), 141[6] (2016), pp. 06016008:1-5.
- 10 Mukhopadhyay, T. and Adhikari, S., "Equivalent in-plane elastic properties of irregular honeycombs: An analytical approach", [International Journal of Solids and Structures](#), 91[8] (2016), pp. 169-184.
- 11 Mukhopadhyay, T. and Adhikari, S., "Effective in-plane elastic properties of auxetic honeycombs with spatial irregularity", [Mechanics of Materials](#), 95[2] (2016), pp. 204-222.

Further reading

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